

Constructing pseudopotentials for ABINIT: about the use of the fhi98PP package

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Description of fhi98PP

1 Generate **norm-conserving** pseudopotentials (DFT) of

- Hamann or Troullier-Martins type
- test their transferability on the atomic level,
including the fully-separable KB form

by a free fortran77 code, operated through
commandline interface using UNIX csh scripts

2 View data by GNU GPL plotting tool xmGrace
<http://plasma-gate.weizmann.ac.il/Grace/>

3 Learn from **tutorial** (ps booklet) and background write-up
M Fuchs, M Scheffler, Comput Phys Commun 119, 67-98 (1999)

→ download from <http://www.FHI-Berlin.MPG.DE/th/fhi98md/fhi98PP/>

Interfaced to ABINIT . . .

. . . gives flexibility where needed

Part I: psatom

Free atom: all-electron full potential \rightarrow pseudo valence orbitals & pseudopotential

- Kohn-Sham equations for full potential \rightarrow eigenstates $\phi_i^{AE}(\mathbf{r}) = \frac{u_{\nu l}^{AE}(r)}{r} Y_{lm}(\Omega)$... spherical

$$\left[-\frac{1}{2} \frac{d^2}{dr^2} + \frac{l(l+1)}{2r^2} + V^{AE}[n^{AE}; r] \right] u_{\nu l}^{AE}(r) = \epsilon_{\nu l}^{AE} u_{\nu l}^{AE}(r), \quad n^{AE}(r) = \sum_{occ} f_i |\phi_i^{AE}(\mathbf{r})|^2$$

▲ Relativity: Dirac \rightarrow scalar relativistic \rightarrow non-relativistic

▲ Full potential $V^{AE}[n^{AE}; r] = -\frac{Z}{r} + V^H[n^{AE}; r] + V^{XC}[n^{AE}; r]$

... XC in LDA or GGA: take same as in solid etc.

- Pseudo atom \rightarrow pseudo valence orbitals $\phi_i(\mathbf{r}) = \frac{u_l(r)}{r} Y_{lm}(\Omega)$ and density n

$$\left[-\frac{1}{2} \frac{d^2}{dr^2} + \frac{l(l+1)}{2r^2} + V_l^{scr}[n; r] \right] u_l(r) = \epsilon_l u_l(r), \quad n(r) = \sum_{occ} f_i |\phi_i(\mathbf{r})|^2$$

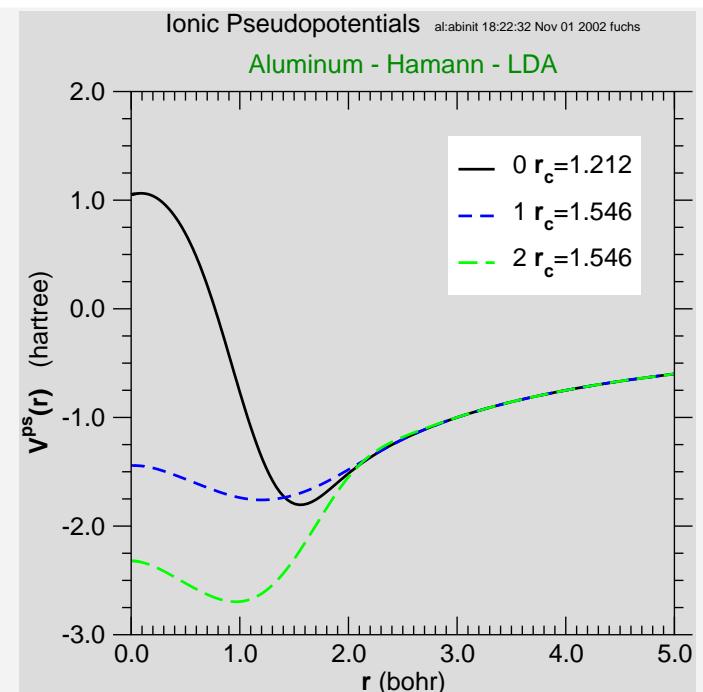
... formally non-relativistic Schrödinger eq.

... different for each valence state $\rightarrow l$ -dependent

→ **Ionic pseudopotentials** $V_l^{ps}[n; r] = V_l^{scr}(r) - V^H[n; r] - V^{XC}[n; r]$

Part I: psatom

Pseudopotential on screen . . .



. . . is **transferable**?

- Start with default cutoff radii & neutral atom:

- Norm-conservation conditions

$$\epsilon_l = \epsilon_{\nu l}^{AE}$$

$$\langle \phi_l | \phi_l \rangle = \langle \phi_{\nu l}^{AE} | \phi_{\nu l}^{AE} \rangle$$

→ ensure correct scattering properties around atomic valence energies

- But how about in the solid or molecule?

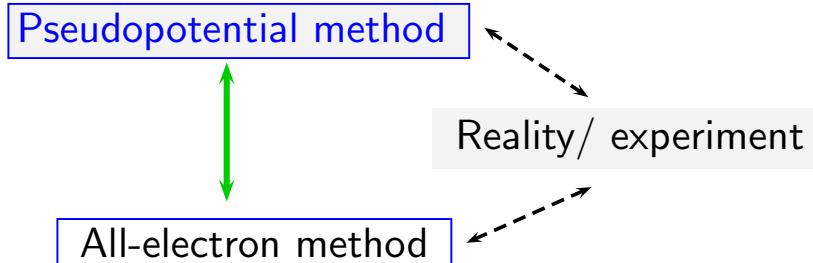
Must test!

. . . and if needed

improved!

- Proposal of how to . . .

Part II: Monitoring transferability with pswatch



- compromise with needed smoothness
- needed accuracy $\sim \mathcal{O}(0.1 \dots 0.01 \text{ eV})$
 - electronic structure
 - cohesive properties
 - atomic structure, relaxation, phonons
 - formation enthalpies, activation energies, ...
- modifications
 - separable potentials (computational)
 - core corrections (methodic)
- ▲ new materials → GaN (with 3d or not), ...
- ▲ new XC functionals → GGA, ...

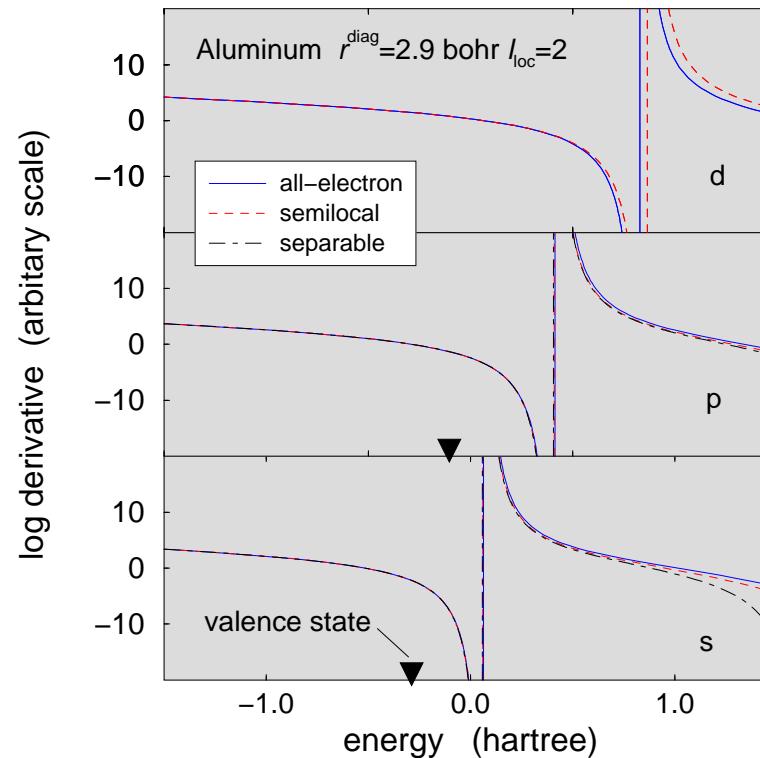
Characteristic tests of PP at atomic level?

Test: Logarithmic derivatives

$$D_l(r^{diag}, \epsilon) = \frac{1}{R_l(\epsilon)} \frac{d}{dr} R_l(r, \epsilon) \Big|_{r^{diag}}$$

norm conservation: o.k. for $\epsilon_l \pm \delta\epsilon$

. . . in practice: over range of valence bands?

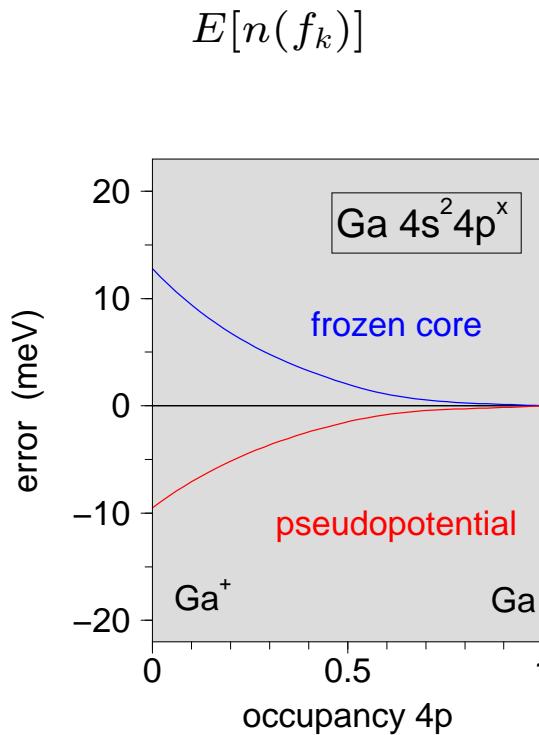


Hardness tests - a „good” example

Test: Configurational changes (ΔSCF)

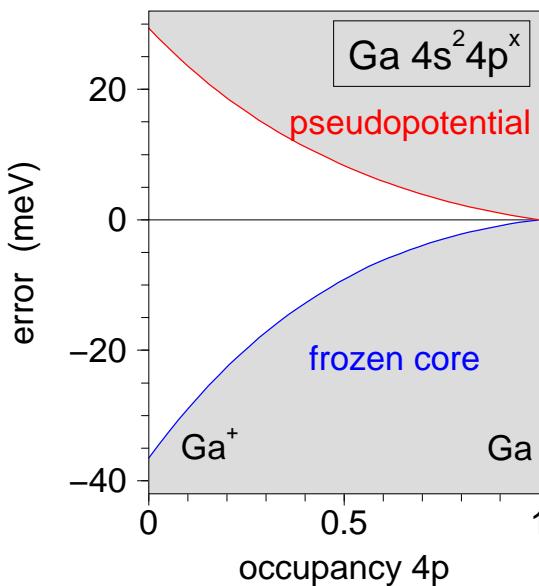
- $s \rightarrow p$ promotion (C, Si, Ge, ...)
- ionization (Li \rightarrow Li $^+$, Na \rightarrow Na $^+$, ...)

- total energy (excitations)



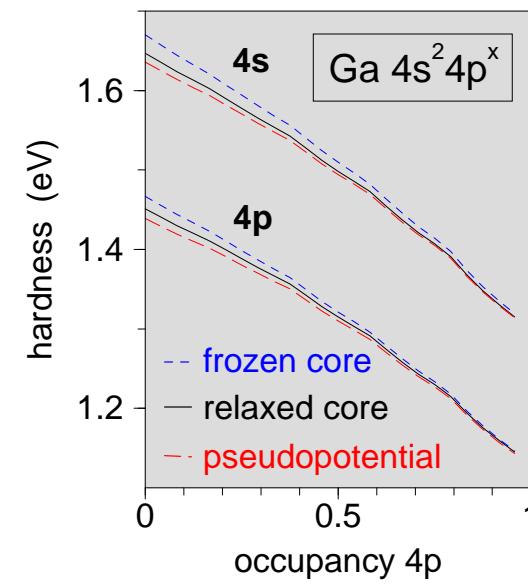
- eigenvalues (Janak theorem)

$$\frac{\partial E(f_k)}{\partial f_i} = \epsilon_i(f_k)$$



- chemical hardness¹

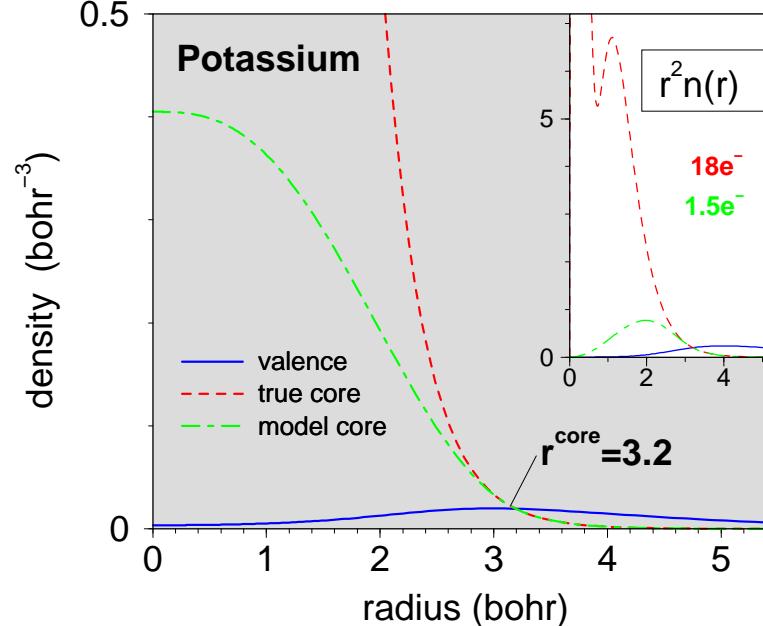
$$\frac{\partial^2 E(f_k)}{\partial f_i \partial f_j} = \frac{\partial \epsilon_i(f_k)}{\partial f_j}$$



¹ Grinberg, Ramer, Rappe, Phys Rev B 63, 201102 (2001); Filipetti et al, Phys Rev B 52, 11793 (1995); Teter, Phys Rev B 48, 5031 (1993).

Nonlinear core-valence XC

- Pseudopotential mimicks also interactions of **valence** and **core** electrons
 - ✓ electrostatic linear in n^v
 - ✗ exchange-correlation **nonlinear**, terms like $(n^c + n^v)^{4/3} \dots$
- so far: linearized core-valence XC
 - ... different in LDA & GGA¹
- restoring nonlinear core-valence XC²
 $E^{XC} = E^{XC}[n^c + n^v]$
 different unscreening:
 $V_l^{ps} = V_l^{scr} - V^H[n^v] - V^{XC}[n^c + n^v]|_{atom}$
- a smooth model core density is used, replaces the true core density inside cutoff radius r^{core}



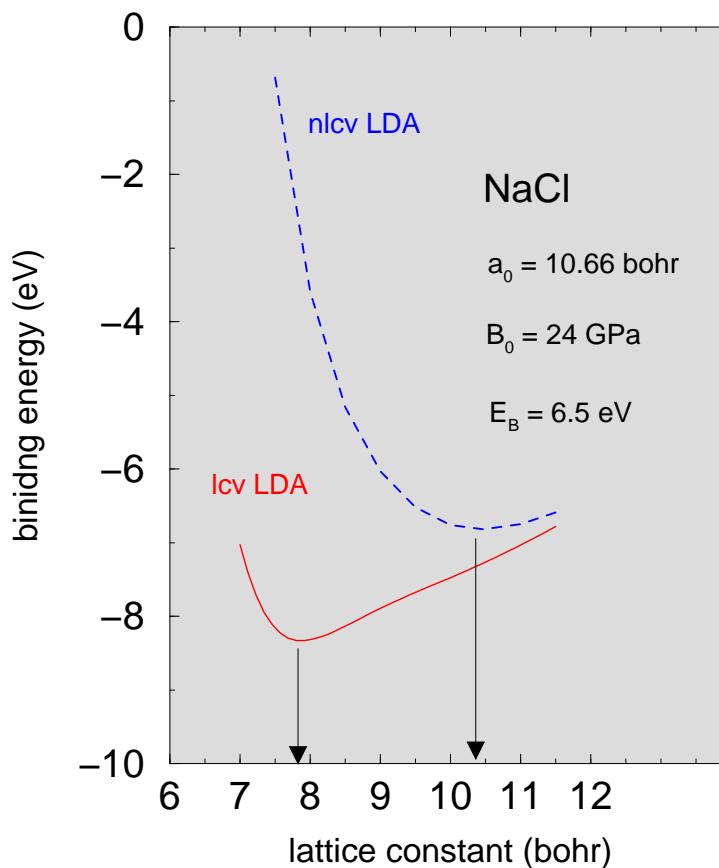
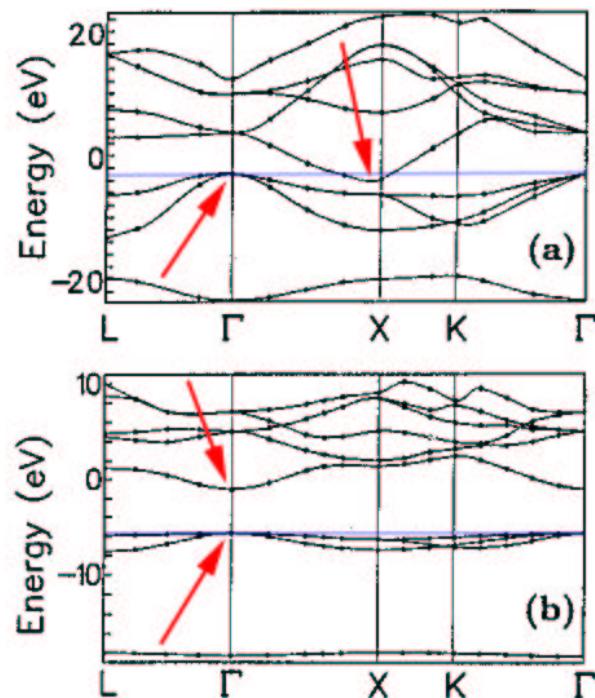
¹ M Fuchs et al, Phys Rev B 57, 2134 (1998)

² SG Louie et al, Phys Rev B 26, 1738 (1982)

... where nonlinear core-valence XC makes a difference

Rocksalt (NaCl):¹

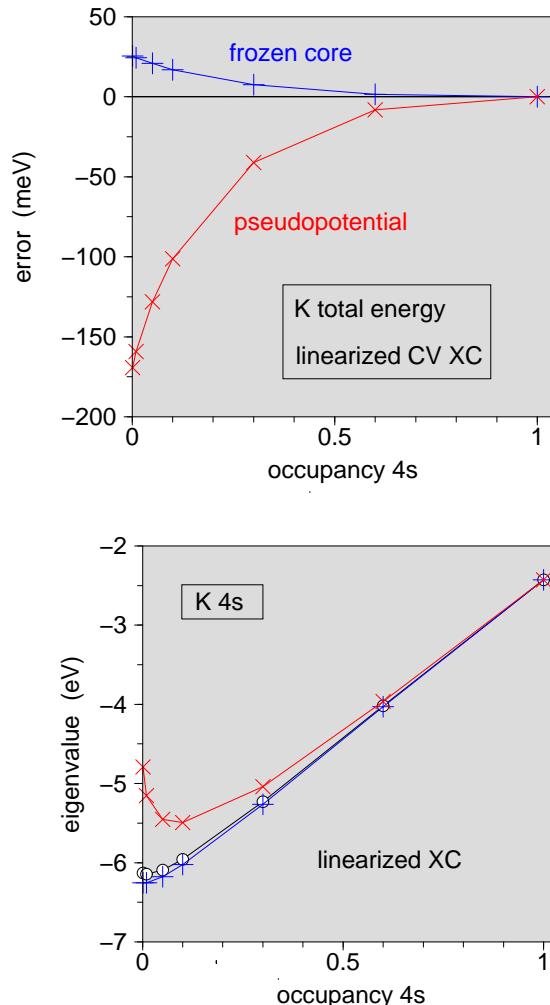
- ✗ semi-metal with linearized CV XC (a)
- ✓ insulator with nonlinear CV XC (b)



¹Hebenstreit, Scheffler, Phys Rev B 46, 10134 (1992).

... and where linearized core-valence XC is fine

Transferability tests would tell. . .



☞ **linearized nlcv XC** mostly sufficient!

- 1st & 2nd row, As, Se, ...
- “two shell” cases → all transition metals,
see Cu: 3-4 XC valence-valence interaction

A test calculation helps...

☞ **nlcv XC** needed:

- “soft” valence shells (alkali’s!)
- extended core states (Zn, Cd, ...) ↔ varying core-valence overlap
- spin-density functional calculations

☞ turning semi-core into valence states?

- Zn 3d, Ca 3d, Rb 4p, ...
- Ga 3d, In 4d in III-nitrides
(but not GaP, GaAs, ...)

... a bit system dependent

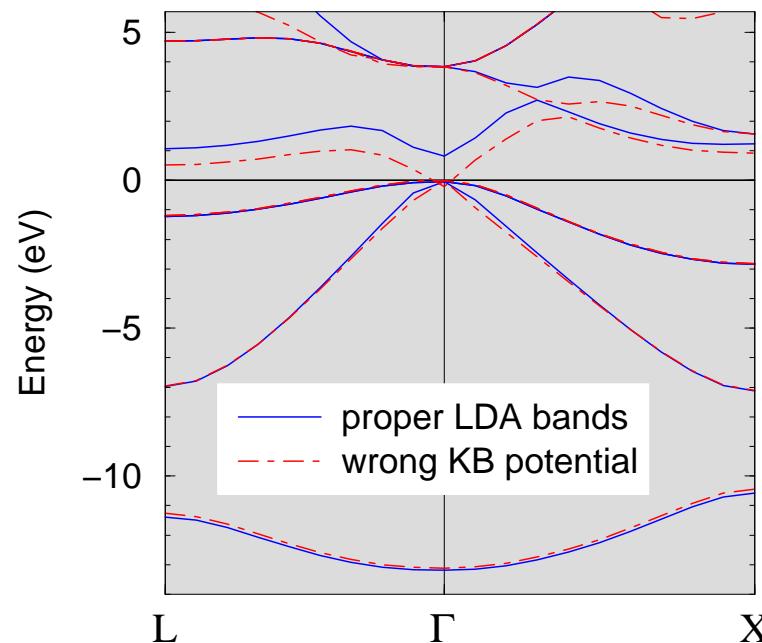
Kleinman-Bylander potentials at work

- ✓ fully separable potentials are computational must
- ☛ full nonlocality → spectral order of states by radial nodes **not** guaranteed
- ➡ **ghost states** above/below physical valence levels possible

- * **Ghost states** detectable in free atom
- ...
- inspect logarithmic derivatives
- do spectral analysis
- ... readily avoided by proper choice of local & nonlocal components

$$V_l(r) \rightarrow \{ V^{loc}(r), \delta V_l(r) \}$$

Example: KB-pseudopotential for As
→ ZB GaAs bandstructure



Analysis of the spectrum of nonlocal Hamiltonians $\hat{H}_l = \hat{T}_l + \hat{V}^{loc} + |\chi_l\rangle\lambda\langle\chi_l|$

- $\lambda = E_l^{KB}$ gives the reference valence level ϵ_l
 - can compare spectra for
 - \square $\lambda = 0$ (local potential only) $\rightarrow \tilde{\epsilon}_i(0)$
 - \square λ arbitrary (with nonlocal potential) $\rightarrow \epsilon_i(\lambda)$

→ for any $\lambda, E_l^{KB} > 0$ spectra ordered like 1

$$\begin{array}{c}
 \lambda \quad : \\
 E_l^{KB} \xrightarrow{\text{dashed}} \tilde{\epsilon}_1(\lambda) \\
 0 \quad \xrightarrow{\text{solid}} \tilde{\epsilon}_1(0) \\
 E_l^{KB} \xrightarrow{\text{solid}} \tilde{\epsilon}_0(\lambda) = \epsilon_l \\
 0 \quad \xrightarrow{\text{solid}} \tilde{\epsilon}_0(0)
 \end{array}$$

no ghost if $\epsilon_l < \tilde{\epsilon}_1(0)$

λ	:	
E_l^{KB}	-----	$\tilde{\epsilon}_2(\lambda)$
0	_____	$\tilde{\epsilon}_2(0)$
E_l^{KB}	_____	$\tilde{\epsilon}_1(\lambda) = \epsilon_l$
0	_____	$\tilde{\epsilon}_1(0)$
E_l^{KB}	_____	$\tilde{\epsilon}_0(\lambda) \times$
0	_____	$\tilde{\epsilon}_0(0)$
ghost if $\epsilon_l > \tilde{\epsilon}_1(0)$		

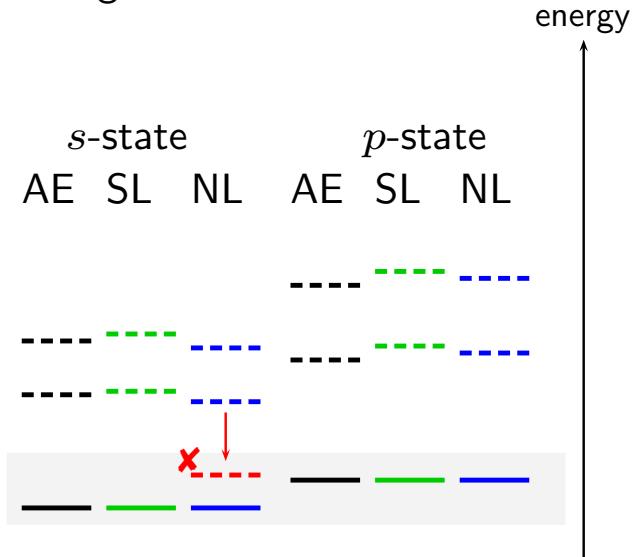
→ for $\lambda, E_l^{KB} < 0$ have $\tilde{\epsilon}_0(\lambda) < \tilde{\epsilon}_0(0) < \tilde{\epsilon}_1(\lambda) \dots$

no ghost if $\epsilon_l < \tilde{\epsilon}_0(0)$

ghost if $\epsilon_l > \tilde{\epsilon}_0(0)$

✿ used as **ghost state criteria** in fhi98PP (pswatch)

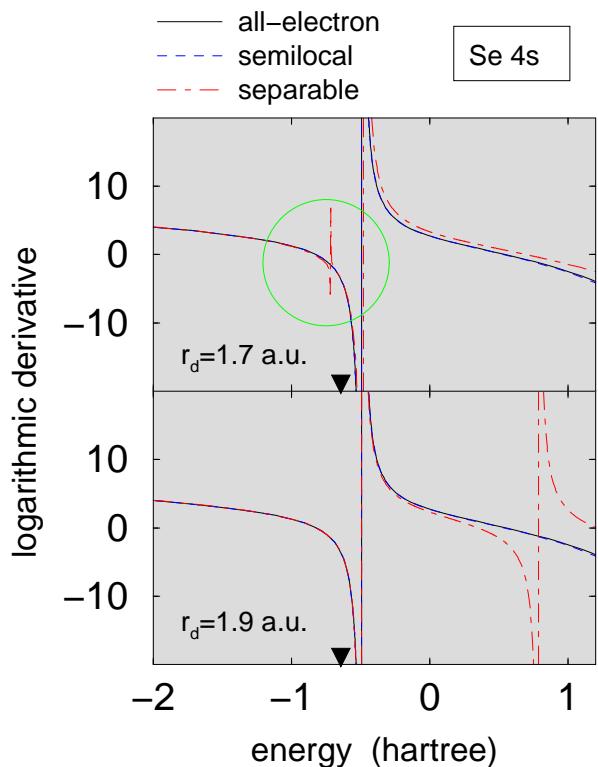
▲ Higher levels o.k. too?



☞ not told by ghost state criteria,
diagonalize $H_l \Rightarrow$ all bound
levels

. . . in case of ghost states

Seen in logarithmic derivatives . . .



. . . where they occur, how to avoid them?

- local potential $l_{loc} = l_{max} = 2$ saves computing
- ✓ unproblematic: 1st & 2nd row, (earth-) alkali's

↙ can cause strong nonlocality (large $|E_l^{KB}|$)

“artificially:” \approx zero denominator in E_l^{KB} (KB-cos)
Ga, Ge, As, Se, ...

vary cutoff radii of local/ nonlocal components

“intrinsically:” numerator of E_l^{KB} large

Cu: deep $V_{3d}(r) \Rightarrow E_l^{KB} \gg 0$ to get 4s right
all 3,4,5d-metals: Cu, Pd, Ag, ...

make local potential repulsive $\rightarrow E_l^{KB} < 0$
use to *s* - or *p*-component!

→ KB-potentials work well in practice

Overview - fhi98PP covers “classic” pseudopotentials

Pseudopotential = electron-i on interaction

- nucleus' Coulomb attraction + core-valence interaction (orthogonality, electrostatic, XC)
 - work throughout periodic table (... almost)
- ✓ physically motivated approximation
 - Valence electrons rule chemical bonding
 - Frozen-core approximation (depends on system)
 - Cancellation of potential and kinetic energy in core
- ✓ well controlled
 - norm-conservation (built in)
 - nonlinear core-valence XC (depends on system)
- ✓ Transferability properties & pseudopotential validation
 - logarithmic derivatives (scattering properties), chemical hardness
 - plane-wave convergence
- ✓ Fully separable, nonlocal potentials
 - analysis & removal of ghost states
- ✓ present XC: LDA & PW91,BP86,PBE,revPBE,RPBE,BLYP GGA . . . easy to add
 - experimental XC: xOEP with KLI approximation, SIC, MGGA
- ✗ desirable: spin-polarization, $L \cdot S$ -coupling,
potential & performance database on the Web ?